

210 B

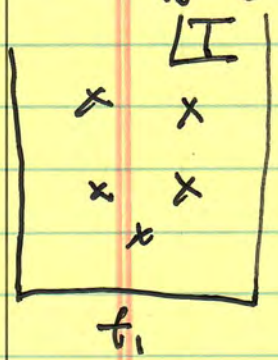
Aggregation - An Application of the Master Equation

(cf: Chandra Review; Krapivsky)
(Schmoluchowski)

Recall Master Eqn:

$$\frac{d P_n}{dt} = \sum_{n'} (W_{n',n} P_{n'} - W_{n,n'}) P_n$$

Now Aggregation: $D \sim T/B$



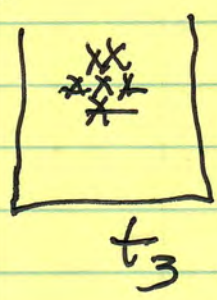
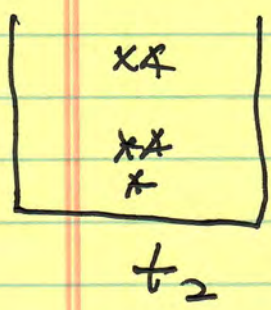
→ system of (charged) colloidal particles, walking randomly

charge colloids in electrolyte

Particles stick when intersect closely



'sticky' collisions



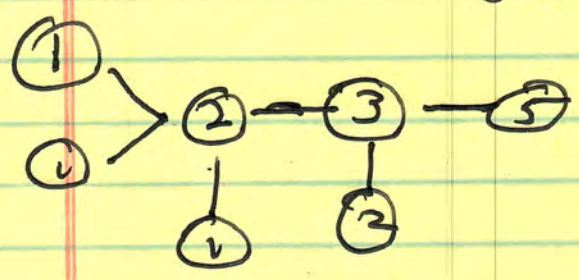
N.B - each particle has 'sphere of influence'. Spheres touch \rightarrow merge

- $R_{50} \sim \lambda D$

- Motion \Rightarrow random walk
 \Rightarrow diffusion

Then, expect:

- evolution
 - $1+1 \rightarrow 2$
 - $2+1 \rightarrow 3$
 - $3+2 \rightarrow 5$



no breakup
 \rightarrow once stuck, remains so.

- we have simultaneous evolution of $P_n \Rightarrow$ probability of n -tuple cluster

- idea then is to model 'birth' and 'death' of n -tuple cluster by interactions of/with other clusters.

so, might have:

$\nu \equiv$ popln. density

$$\frac{d}{dt} \nu_2 = \underbrace{c\nu}_{1+4} - \underbrace{\text{out}}_{2+1, 3, 4, \dots}$$

$$\frac{d}{dt} \nu_3 = \underbrace{c\nu}_{2+4} - \underbrace{\text{out}}_{3+1, 2, \dots}$$

etc. \leftarrow to ν_4

in general,

rate production of n-particle

$$\approx \sum_{p, z} \left(\# \underbrace{\nu_p \nu_z}_{p+z=n} \right)$$

rate coeff \swarrow

joint popln / prob. $p+z$

rate destruction of n-particle

$$\approx \sum_{i=1}^N \# \nu_n \nu_i$$

i.e. n-particle + any other.

So, have:

- classic input-output / birth and death model for populations

- system of coupled Master Eqns

i.e.

birth \rightarrow merger of smaller

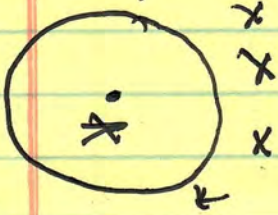
$$\frac{d}{dt} v_n = \sum_{\substack{i+j \\ i+j=n}} C_{ij} v_i v_j - \sum_{l=1}^n C_{ln} v_l v_n$$

death \rightarrow absorption to larger

What are the rate coefficients?

\rightarrow Rate of interaction \Rightarrow diffusion

\rightarrow density, diffuses = consider rate at which sphere R sweeps fluid



$$n = \text{const}, \quad t = 0$$

$$|v| > R$$

\uparrow
 R_{si}

$$\frac{\partial n}{\partial t} = D \nabla^2 n$$

$$n=0, \quad r=R \quad t \rightarrow \infty$$

absorbing boundary

Heuristically, since particles diffuse together

$$\frac{\partial n}{\partial t} \sim 4\pi R^2 \left(+D \frac{dn}{dr} \Big|_R \right)$$

Area of Influence Sphere Flux thru surface of sphere

→

$$\frac{\partial n}{\partial t} \sim 4\pi D R^2 D \frac{dn}{dr}$$

$$\sim 4\pi D R^2 \frac{D}{R} \sim 4\pi D R \frac{D}{C_{ij}}$$

N.B.

$$\frac{\partial n}{\partial t} = D \nabla^2 n$$

symmetry →

$$Q_+(n) = D \frac{d^2}{dn^2} (n) t$$

$$x = nr, \quad x|_{t \rightarrow \infty} = a + br$$

$$n = r \left[1 - \frac{R}{r} + \frac{2R}{r\sqrt{\pi}} \int_0^{(n-R)/(2Dt)^{1/2}} e^{-x^2} dx \right]$$

18

$$\text{Rate}_{\text{swept}} = 4\pi DR r (1 + R) (\pi Dt)^{1/2}$$

Now, Rate for mergers:

i sweeps k

$$J_{i,k} dt = 4\pi D_{i,k} v_i v_k dt$$

$$D_{i,k} = D_i + D_k$$

→ independent diffusion

→ mean square step adds (quadratures)

19

to avoid double counting

$$\frac{d}{dt} v_k = 4\pi \left(\begin{array}{l} \text{in} \\ \frac{1}{2} \sum_{\substack{i,j \\ i+j=k}} v_i v_j D_{j,i} R_{ij} \\ \text{out} \\ - v_k \sum_{j=1}^k v_j D_{k,j} R_{k,j} \end{array} \right)$$

- Master Eqn. for k -tuple
(Really N)

Now, simplifying approximations:

- $R_i = R_k = R$ single scale
- $D_j = D$ all same
- $D_i R_i = DR$ "
- $D_{in} R_{in} = 2DR$ D 's edge

So have much simpler :

$$\frac{d}{dt} v_k = 8\pi DR \left(\frac{4}{2} \sum_{\substack{i+j \\ =k}} v_i v_j - v_k \sum_{j=1}^{\infty} v_j \right)$$

rate

$$\tau = 4\pi DR t$$

$$\frac{dv_k}{d\tau} = \sum_{i+j} v_i v_j - 2v_k \sum_{j=1}^{\infty} v_j$$

→ reduced Master Equation

→ specifies population evolution for clusters.

then, to solve:

$$\sum_k$$

→

$$\begin{aligned} \frac{d}{dt} \left(\sum_{k=1}^{\infty} v_k \right) &= \sum_i \sum_j v_i v_j - 2 \sum_{k=1}^{\infty} \sum_{j=1}^{\infty} v_k v_j \\ &\quad \text{reless on cluster} \\ &= - \left(\sum_{k=1}^{\infty} v_k \right)^2 \end{aligned}$$

$$\frac{dx}{dt} = -x^2 \quad \rightarrow \quad 1/x = t + C$$

$$\sum_{k=1}^{\infty} v_k = v_0 / (1 + \tau v_0)$$

sum of
populations decays
 \rightarrow cluster into 1.

And can solve for populations:

$$\begin{aligned} \frac{d}{dt} v_i &= -2v_i \sum_{k=1}^{\infty} v_k \\ &= -2v_i \sum v_k \\ &= -2v_i v_0 / (1 + \tau v_0) \end{aligned}$$

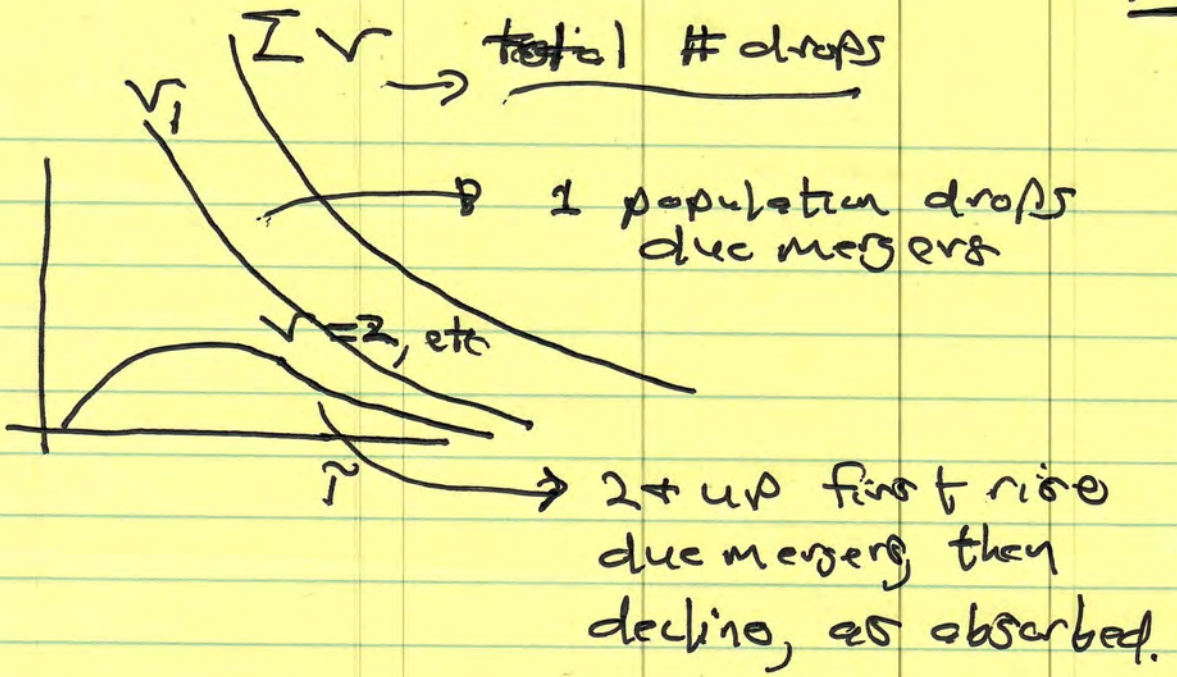
$$v_i = v_0 / (1 + \tau v_0)^2$$

and

$$v_k = v_0 (\tau v_0)^{k-1} / (1 + \tau v_0)^{k+1}$$

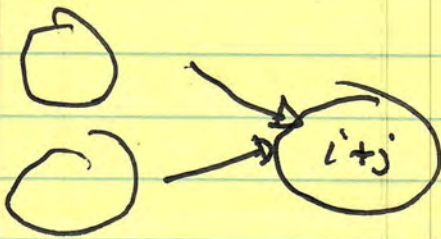
$$k = 1, 2, 3, \dots$$

50



N.B:

- typical of aggregation problems



clusters join (irreversibly) when meet

eg. { milk curdling
blood coagulation
planet formation

[1 way flow → larger scale]

- basic description is system of Master Eqs

$$\frac{d}{dt} C_k = \frac{4}{2} \sum_{i+j=k} \overset{\text{in}}{k_{ij}} c_i c_j - c_k \sum_{l \geq 1} \overset{\text{out}}{k_{kl}} C_l$$

couplings are the key!
 \Rightarrow rate

\Rightarrow time to form large (∞)
 cluster finite, infinite
 [finite time singularity?]

- Typical of 'inverse cascade'.

- Many approaches to solution

- exact \rightarrow some simple cases

- moments

- recursion

see Krepiuski

→ Constraint: "Basic Equation conserves mass."

$$C \rightarrow F$$

i.e. nothing lost enroute to N cluster.

$$M \rightarrow \sum k c$$

$$M(t) = \sum_{k \geq 1} k c_k(t)$$

Now

$$\frac{dM}{dt} = \sum_k k \frac{dc_k}{dt}$$

$$= \sum_k \sum_{\substack{i, j \\ i+j=k}} \frac{1}{2} k_{ij} c_i c_j$$

$$- \sum_k \sum_i k_{ik} c_i c_k$$

$$i+j=k$$

$$= \sum_k \sum_i \frac{1}{2} (2) k_{i,k} c_i c_k$$

$$- \sum_k \sum_i k_{ik} c_i c_k$$

$$= 0$$

so, Mass conserved, (plausible).

this brings us to Gelation

⇒

Gelation = Aggregation with rate
increasing with cluster mass
(i.e. more interesting 'coeffs').