

**PHYSICS 239.c : CONDENSED MATTER PHYSICS
FINAL EXAMINATION**

Instructions: Do problem 0 and any three of problems 1 through 4.

(0) Provide brief but accurate answers to each of the following questions:

- (a) What is the Hohenberg-Mermin-Wagner theorem, and what is Goldstone's theorem?
- (b) In the context of the Boltzmann equation, what is meant by the term, "collisional invariant"? What are two examples of collisional invariants in the case of (single band) electron transport?
- (c) The point group D_8 , describing the symmetries of a planar octagon, is relevant to molecular chemistry, but is not among the 32 crystallographic point groups. Why not?
- (d) What is the Mössbauer effect?
- (e) What is a Wannier state? What quantum numbers are necessary to specify a Wannier state? What completeness and orthonormality conditions to the Wannier states satisfy?

(1) The hexagonal close packed (hcp) structure is a simple hexagonal (sh) Bravais lattice with a two-element basis. The three elementary direct lattice vectors of the sh structure are

$$\mathbf{a}_1 = a\left(\frac{1}{2}\hat{x} - \frac{\sqrt{3}}{2}\hat{y}\right) \quad , \quad \mathbf{a}_2 = a\left(\frac{1}{2}\hat{x} + \frac{\sqrt{3}}{2}\hat{y}\right) \quad , \quad \mathbf{a}_3 = c\hat{z}$$

The two basis vectors are $\mathbf{0}$ and $\boldsymbol{\delta} = \frac{1}{3}\mathbf{a}_1 + \frac{2}{3}\mathbf{a}_2 + \frac{1}{2}\mathbf{a}_3$.

- (a) The hcp lattice is close-packed, which means that $|\boldsymbol{\delta}| = a$. Find the value of c in terms of the in-plane lattice spacing a .
- (b) What is the coordination number z (*i.e.* the number of nearest neighbors of any given site) of the hcp lattice? Write down the positions of all z neighbors of the lattice site $\mathbf{0}$ in terms of $\mathbf{a}_{1,2,3}$ and $\boldsymbol{\delta}$.
- (c) What are the three elementary reciprocal lattice vectors \mathbf{b}_1 , \mathbf{b}_2 , and \mathbf{b}_3 ?
- (d) The space group of the hcp structure ($P6_3/mmc$) is nonsymmorphic (it contains a twofold screw operation). Consider the Bragg peaks located at wavevectors $\mathbf{G} = n_1\mathbf{b}_1 + n_2\mathbf{b}_2 + n_3\mathbf{b}_3$. What is the condition on $\{n_1, n_2, n_3\}$ for there to be an extinction in the diffraction pattern at \mathbf{G} ?

(2) Consider the tight binding Hamiltonian for s -orbitals on the railroad trestle lattice, depicted in Fig. 1. The hopping amplitude along each rail is t and the hopping amplitude between the rails is t' . Both t and t' are positive.

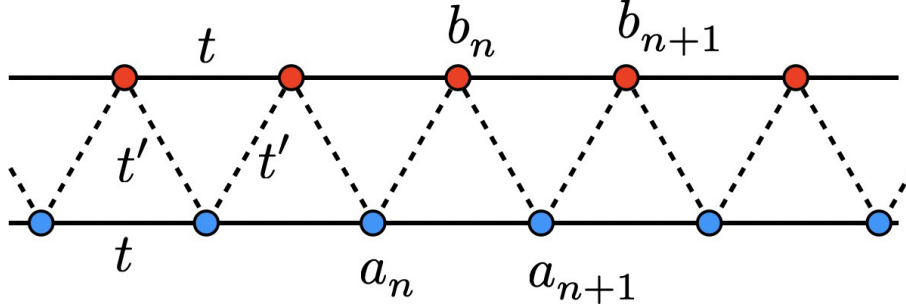


Figure 1: Railroad trestle lattice

- (a) Write the tight-binding Hamiltonian in real space. You may use either bracket notation with local orthonormal orbitals $|n, a\rangle$ and $|n, b\rangle$ or fermionic second quantized operators a_n and b_n and their conjugates.
- (b) Write the tight-binding Hamiltonian in crystal momentum space, *i.e.* using the Fourier transformed states $|k, a\rangle$ and $|k, b\rangle$ or the second quantized operators a_k and b_k (and their conjugates).
- (c) Solve for the electronic energy bands $E_j(\theta)$, where $\theta = ka$ and a is the lattice spacing along either rail. How many bands are there? Sketch their dispersion.

(3) Consider an infinite one-dimensional chain of atoms, each of mass m , located at positions $x_n = na + u_n$, with potential energy

$$V = \frac{1}{2} \sum_{n < n'} K_{nn'} (u_n - u_{n'})^2 \quad .$$

Thus, each pair of atoms (n, n') is connected by a spring of spring constant $K_{nn'}$ whose unstretched length is $|n - n'|a$, where a is the lattice constant. You may assume the potential has the lattice translation symmetry, *i.e.* $K_{nn'} = K(n - n') = K(n' - n)$ is an even function of the difference $n - n'$.

- (a) Find the equation of motion for the Fourier modes $\hat{u}_k \equiv N^{-1/2} \sum_n u_n e^{-ikna}$, where $N \rightarrow \infty$ is the number of unit cells.
- (b) Find an expression for the phonon dispersion $\omega(k)$.
- (c) Write down an expression for the ground state wavefunction $\Psi_0(\{u_n\})$.
- (d) Suppose $K(\ell) = K_0 \ell^{-2}$. Compute the phonon frequency $\omega(k)$ and the zero temperature quantum fluctuation $\langle \Psi_0 | u_n^2 | \Psi_0 \rangle$ of the atomic positions. It may interest you to know that for $\theta \in [0, 2\pi]$, it is a True Fact that

$$\text{Re Li}_2(e^{i\theta}) = \sum_{n=1}^{\infty} \frac{\cos(n\theta)}{n^2} = \frac{1}{6}\pi^2 - \frac{1}{4}\theta(2\pi - \theta) \quad ,$$

where

$$\text{Li}_k(z) = \sum_{n=1}^{\infty} \frac{z^n}{n^k} .$$

is the polylogarithm function.

(e) For $K(\ell) = K_0 (\delta_{\ell,1} + \delta_{\ell,-1})$, we found $\omega_k = 2(K_0/m)^{1/2} |\sin(\frac{1}{2}ka)|$, hence $\omega(k) = c|k|$ at long wavelengths. The zero temperature fluctuations $\Psi_0(\{u_n\})$ then diverge. Yet your result for the fluctuations in part (d) should have been finite. Why do you suppose this might be the case?

(4) Consider the currents

$$\begin{aligned} \mathbf{j} &= -2e \int_{\hat{\Omega}} \frac{d^3k}{(2\pi)^3} \mathbf{v} \delta f \\ \mathbf{J} &\equiv 2 \int_{\hat{\Omega}} \frac{d^3k}{(2\pi)^3} (\varepsilon - \mu)^2 \mathbf{v} \delta f . \end{aligned}$$

Define the response coefficients ρ , Q , ω , and v by the relations

$$\begin{aligned} \mathcal{E} &= \rho \mathbf{j} + Q \nabla T \\ \mathbf{J} &= \omega \mathbf{j} - v \nabla T . \end{aligned}$$

For a system with cubic symmetry, find expressions for the transport coefficients ρ , Q , ω , and v in terms of the integrals

$$\begin{aligned} \mathcal{K}_n &= \frac{\tau}{12\pi^3 \hbar} \int_{-\infty}^{\infty} d\varepsilon (\varepsilon - \mu)^n \left(-\frac{\partial f^0}{\partial \varepsilon} \right) \int dS_{\varepsilon} |\mathbf{v}| \\ &= \frac{\sigma_0}{e^2} \varepsilon_F^{-3/2} \mathcal{S} [\varepsilon^{3/2} (\varepsilon - \mu)^n] \Big|_{\varepsilon=\mu} , \end{aligned}$$

where

$$\mathcal{S} = \pi \mathcal{D} \csc \pi \mathcal{D} = 1 + \frac{\pi^2}{6} \mathcal{D}^2 + \frac{7\pi^4}{360} \mathcal{D}^4 + \dots ,$$

with $\mathcal{D} = k_B T \partial_{\varepsilon}$.