PHYSICS 152B/232 Spring 2017 Homework Assignment #3 Solutions

[1] Cyclotron resonance in Si and Ge – Both Si and Ge are indirect gap semiconductors with anisotropic conduction band minima and doubly degenerate valence band maxima. In Si, the conduction band minima occur along the $\langle 100 \rangle$ ($\langle \Gamma X \rangle$) directions, and are six-fold degenerate. The equal energy surfaces are cigar-shaped, and the effective mass along the $\langle \Gamma X \rangle$ principal axes (the 'longitudinal' effective mass) is $m_1^* \simeq 1.0 m_e$, while the effective mass in the plane perpendicular to this axis (the 'transverse' effective mass) is $m_t^* \simeq 0.20 m_e$. The valence band maximum occurs at the unique Γ point, and there are two isotropic hole branches: a 'heavy' hole with $m_{\rm hh}^* \simeq 0.49 \, m_{\rm e}$, and a 'light' hole with $m_{\rm lh}^* \simeq 0.16 \, m_{\rm e}$.

In Ge, the conduction band minima occur at the fourfold degenerate L point (along the eight $\langle 111 \rangle$ directions) with effective masses $m_1^* \simeq 1.6 m_e$ and $m_t^* \simeq 0.08 m_e$. The valence band maximum again occurs at the Γ point, where the hole masses are $m_{\text{hh}}^* \simeq 0.34 m_{\text{e}}$ and $m_{\text{lh}}^* \simeq 0.044 m_{\text{e}}$. Use the following figures to interpret the cyclotron resonance data shown below. Verify whether the data corroborate the quoted values of the effective masses in Si and Ge.

Solution :

We found that $\sigma_{\alpha\beta} = ne^2 \Gamma_{\alpha\beta}^{-1}$, with

$$
\Gamma_{\alpha\beta} \equiv (\tau^{-1} - i\omega) m_{\alpha\beta} \pm \frac{e}{c} \epsilon_{\alpha\beta\gamma} B^{\gamma}
$$
\n
$$
= \begin{pmatrix}\n(\tau^{-1} - i\omega) m_x^* & \pm eB_z/c & \mp eB_y/c \\
\mp eB_z/c & (\tau^{-1} - i\omega) m_y^* & \pm eB_x/c \\
\pm eB_y/c & \mp eB_x/c & (\tau^{-1} - i\omega) m_z^*\n\end{pmatrix}.
$$

The valence band maxima are isotropic in both cases, with

$$
m_{\rm hh}^{*}(\rm Si) \simeq 0.49 m_{\rm e}
$$
 $m_{\rm hh}^{*}(\rm Ge) \simeq 0.34 m_{\rm e}$
 $m_{\rm lh}^{*}(\rm Si) \simeq 0.16 m_{\rm e}$ $m_{\rm lh}^{*}(\rm Ge) \simeq 0.044 m_{\rm e}$.

With isotropic bands, the absorption is peaked at $\omega = \omega_c = eB/m^*c$, assuming

Figure 1: Constant energy surfaces near the conduction band minima in silicon. There are six symmetry-related ellipsoidal pockets whose long axes run along the $\langle 100 \rangle$ directions.

 $\omega_{\rm c}\tau \gg 1$. Writing $\omega = 2\pi f$, the resonance occurs at a field

$$
B(f) = 2\pi f \cdot \frac{m^* c}{e}
$$

= $\frac{hc}{e} \cdot \frac{m^*}{m_e} \cdot \frac{1}{2\pi a_B^2} \cdot \frac{hf}{(e^2/a_B)}$
= $3.58 \times 10^{-7} \text{ G} \cdot \frac{m^*}{m_e} \cdot f[\text{Hz}]$
= $8590 \text{ G} \cdot \frac{m^*}{m_e}$,

Figure 2: Cyclotron resonance data in Si (G. Dresselhaus et al., Phys, Rev, 98, 368 (1955).) The field lies in a (110) plane and makes an angle of $30°$ with the [001] axis.

where we have used

$$
\frac{hc}{e} = 4.137 \times 10^{-7} \,\text{G} \cdot \text{cm}^2
$$

$$
a_{\text{B}} = \frac{\hbar^2}{m_{\text{e}}e^2} = 0.529 \,\text{\AA}
$$

$$
h = 4.136 \times 10^{-15} \,\text{eV} \cdot \text{s}
$$

$$
\frac{e^2}{a_{\text{B}}} = 27.2 \,\text{eV} = 2 \,\text{Ry}
$$

$$
f = 2.40 \times 10^{10} \,\text{Hz}
$$

Thus, we predict

$$
B_{hh}(Si) \simeq 4210 \text{ G}
$$
 $B_{hh}(Ge) \simeq 2920 \text{ G}$
\n $B_{lh}(Si) \simeq 1370 \text{ G}$ $B_{lh}(Ge) \simeq 378 \text{ G}$.

All of these look pretty good.

Figure 3: Constant energy surfaces near the conduction band minima in germanium. There are eight symmetry-related half-ellipsoids whose long axes run along the $\langle 111 \rangle$ directions, and are centered on the midpoints of the hexagonal zone faces. With a suitable choice of primitive cell in k-space, these can be represented as four ellipsoids, the half-ellipsoids on opposite faces being joined together by translations through suitable reciprocal lattice vectors.

Now let us review the situation with electrons near the conduction band minima:

Si: 6-fold degenerate minima along $\langle 100 \rangle$ Ge: 4-fold degenerate minima along $\langle 111 \rangle$ (at L point)

The resonance condition is that $\sigma_{\alpha\beta} = \infty$, which for $\tau > 0$ occurs only at complex frequencies, *i.e.* for real frequencies there are no true divergences, only resonances. The location of the resonance is determined by det $\Gamma = 0$. Taking the determinant, one finds

$$
\det \Gamma = (\tau^{-1} - i\omega) m_1^* \cdot \left\{ (\tau^{-1} - i\omega)^2 m_t^{*2} + \frac{e^2}{c^2} B_z^2 + \frac{m_t^*}{m_1^*} \frac{e^2}{c^2} \left(B_x^2 + B_y^2 \right) \right\} .
$$

Figure 4: Cyclotron resonance data in Ge (G. Dresselhaus et al., Phys, Rev, 98, 368 (1955).) The field lies in a (110) plane and makes an angle of $60°$ with the [001] axis.

Assuming $\omega \tau \gg 1$, the location of the resonance is given by

$$
\omega^2 = \left(\frac{eB_\parallel}{m_{\rm t}^*c}\right)^2 + \frac{m_{\rm t}^*}{m_{\rm l}^*}\left(\frac{e{\pmb B}_\perp}{m_{\rm t}^*c}\right)^2\,, \label{eq:omega2}
$$

where $B_{\parallel} \equiv B_z$ and $\mathbf{B}_{\perp} \equiv B_x \hat{\mathbf{x}} + B_y \hat{\mathbf{y}}$. Let the polar angle of \mathbf{B} be θ , so $B_{\parallel} = B \cos \theta$ and $B_{\perp} = B \sin \theta$. We then have

$$
\omega^2 = \left(\frac{eB}{m_{\rm t}^*c}\right)^2 \left\{ \cos^2 \theta + \frac{m_{\rm t}^*}{m_{\rm l}^*} \sin^2 \theta \right\}
$$

$$
B(f) = 8600 \,\mathrm{G} \cdot \left(\frac{m_{\rm t}^*}{m_{\rm e}}\right) / \sqrt{\cos^2 \theta + \frac{m_{\rm t}^*}{m_{\rm l}^*} \sin^2 \theta} \;,
$$

where again we take $f = \omega/2\pi = 2.4 \times 10^{10}$ Hz.

According to the diagrams, the field lies in the (110) plane, which means we can write

$$
\hat{\boldsymbol{B}} = \sqrt{\frac{1}{2}} \sin \chi \,\hat{\boldsymbol{e}}_1 - \sqrt{\frac{1}{2}} \sin \chi \,\hat{\boldsymbol{e}}_2 + \cos \chi \,\hat{\boldsymbol{e}}_3 \ ,
$$

where χ is the angle $\hat{\mathbf{B}}$ makes with $\hat{\mathbf{e}}_3 = [001]$.

Ge : We have

$$
\frac{m_{\rm t}^*}{m_{\rm e}} = 0.082 \qquad \frac{m_{\rm t}^*}{m_{\rm l}^*} = 0.051 \; , \label{eq:mass}
$$

and we are told $\chi = 60^{\circ}$, so

$$
\hat{B} = \sqrt{\frac{3}{8}} \hat{e}_1 - \sqrt{\frac{3}{8}} \hat{e}_2 + \frac{1}{2} \hat{e}_3.
$$

The conduction band minima lie along $\langle 111 \rangle$, which denotes a set of directions in real space:

$$
\pm[111]: \quad \hat{n} = \pm \frac{1}{\sqrt{3}}(\hat{e}_1 + \hat{e}_2 + \hat{e}_3) \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{1}{12} \quad \Rightarrow \quad B = 1950 \text{ G}
$$
\n
$$
\pm[11\bar{1}]: \quad \hat{n} = \pm \frac{1}{\sqrt{3}}(\hat{e}_1 + \hat{e}_2 - \hat{e}_3) \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{1}{12} \quad \Rightarrow \quad B = 1950 \text{ G}
$$
\n
$$
\pm[\bar{1}11]: \quad \hat{n} = \pm \frac{1}{\sqrt{3}}(-\hat{e}_1 + \hat{e}_2 + \hat{e}_3) \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{7-2\sqrt{6}}{12} \quad \Rightarrow \quad B = 1510 \text{ G}
$$
\n
$$
\pm[1\bar{1}1]: \quad \hat{n} = \pm \frac{1}{\sqrt{3}}(\hat{e}_1 - \hat{e}_2 + \hat{e}_3) \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{7+2\sqrt{6}}{12} \quad \Rightarrow \quad B = 710 \text{ G}
$$
\nAll OK!

 Si : Again, *B* lies in the (110) plane, this time with $\chi = 30^{\circ}$, so

$$
\hat{B} = \sqrt{\frac{1}{8}} \,\hat{e}_1 - \sqrt{\frac{1}{8}} \,\hat{e}_2 + \sqrt{\frac{3}{4}} \,\hat{e}_3 \ .
$$

The conduction band minima lie along $\langle 100 \rangle$, so

$$
\pm[001]: \quad \hat{n} = \pm \hat{e}_3 \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{3}{4} \quad \Rightarrow \quad B = 1820 \text{ G}
$$
\n
$$
\pm[010]: \quad \hat{n} = \pm \hat{e}_2 \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{1}{8} \quad \Rightarrow \quad B = 2980 \text{ G}
$$
\n
$$
\pm[100]: \quad \hat{n} = \pm \hat{e}_1 \quad \Rightarrow \quad \cos^2 \theta = (\hat{B} \cdot \hat{n})^2 = \frac{1}{8} \quad \Rightarrow \quad B = 2980 \text{ G}
$$

These also look pretty good.