

$$8-11. \quad \frac{n_2}{n_1} = \frac{g_2 e^{-E_2/kT}}{g_1 e^{-E_1/kT}} = \frac{g_2}{g_1} e^{-E_2-E_1/kT}$$

$$e^{E_2-E_1/kT} = \frac{g_2}{g_1} \times \frac{n_1}{n_2} = E_2 - E_1 / kT = \ln \left(\frac{g_2}{g_1} \times \frac{n_1}{n_2} \right)$$

$$T = \frac{E_2 - E_1}{k \ln \left[\frac{g_2}{g_1} \frac{n_1}{n_2} \right]} = \frac{10.2 \text{ eV}}{8.617 \times 10^{-5} \text{ eV/K} \ln 4 \times 10^6} = 7790 \text{ K}$$

$$8-12. \quad \frac{n_2}{n_1} = \frac{g_2}{g_1} e^{-E_2-E_1/kT} = \frac{3}{1} e^{-\left[\frac{4 \times 10^{-3} \text{ eV}}{8.617 \times 10^{-5} \text{ eV/K} \cdot 300 \text{ K}} \right]} = 2.57$$

8-17. For hydrogen: $E_n = -\frac{mk^2e^4}{2\hbar^2} \frac{1}{n^2} = -\frac{13.605687}{n^2} eV$ using values of the constants accurate to six decimal places.

$$E_1 = -13.605687 eV$$

$$E_2 = -3.401422 eV \quad E_2 - E_1 = 10.204265 eV$$

$$E_3 = -1.511743 eV \quad E_3 - E_1 = 12.093944 eV$$

$$(a) \frac{n_2}{n_1} = \frac{g_2}{g_1} e^{-E_2 - E_1 / kT} = \frac{8}{2} e^{-10.20427 / 0.02586} = 4e^{-395} = 4 \times 10^{-172} \approx 0$$

$$\frac{n_3}{n_1} = \frac{g_3}{g_1} e^{-E_3 - E_1 / kT} = \frac{18}{2} e^{-12.09394 / 0.02586} = 9e^{-468} = 9 \times 10^{-203} \approx 0$$

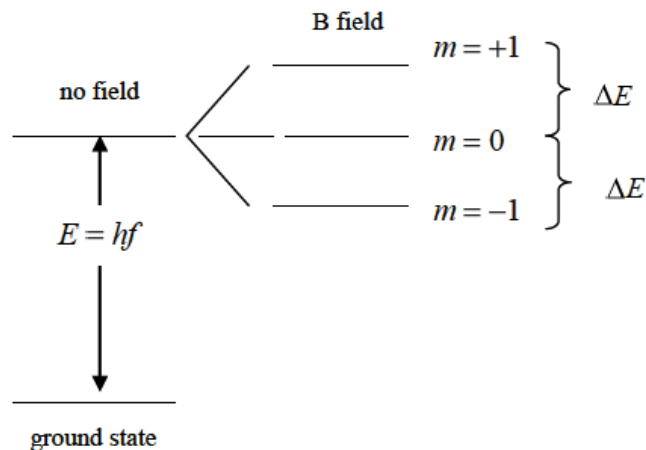
$$(b) \frac{n_2}{n_1} = 0.01 = 4e^{-10.20427/kT} \rightarrow e^{-10.20427/kT} = 0.0025$$

$$-10.20427/kT = \ln 0.0025 = -5.99146$$

$$T = \frac{10.20427 eV}{5.99146 \cdot 8.61734 \times 10^{-5} eV \cdot K} = 19,760 K$$

$$(c) \frac{n_3}{n_1} = 9e^{-12.09394 / 8.61734 \cdot 10^{-5} \cdot 19,760} = 0.00742 = 0.7\%$$

8-18.



Neglecting the spin, the $3p$ state is doubly degenerate: $\ell = 0, 1$ hence, there are two $m = 0$ levels equally populated.

$$E = hf = hc / \lambda = 1.8509 eV \quad \lambda = 670.79 nm$$

$$\Delta E = \frac{ehB}{2m_e} = 2.315 \times 10^{-4} eV$$

(a) The fraction of atoms in each m -state relative to the ground state is: (Example 8-2)

$$\frac{n_{+1}}{n} = e^{-1.8511/0.02586} = e^{-71.58} = 10^{-31.09} = 8.18 \times 10^{-32}$$

$$\frac{n_0}{n} = 2 \times e^{-1.8509/0.02586} = 2e^{-71.57} = 2 \times 10^{-31.08} = 1.64 \times 10^{-31}$$

$$\frac{n_0}{n_{-1}} = e^{-1.8507/0.02586} = e^{-71.56} = 10^{-31.08} = 8.30 \times 10^{-32}$$

(b) The brightest line with the B-field “on” will be the transition from the $m = 0$ level, the center line of the Zeeman spectrum. With that as the “standard”, the relative intensities will be: $8.30/16.4/8.18 \rightarrow 0.51/1.00/0.50$

8-21. Assuming the gasses are ideal gases, the pressure is given by: $P = \frac{2}{3} \frac{N \langle E \rangle}{V}$ for classical, FD, and BE particles. P_{FD} will be highest due to the exclusion principle, which, in effect, limits the volume available to each particle so that each strikes the walls more frequently than the classical particles. On the other hand, P_{BE} will be lowest, because the particles tend to be in the same state, which in effect, is like classical particles with a mutual attraction, so they strike the walls less frequently.

$$8-23. \quad \lambda = \frac{h}{p} = \frac{h}{\sqrt{2m \langle E \rangle}} = \frac{h}{\sqrt{2m \cdot 3kT/2}} = \frac{h}{3mkT^{1/2}}$$

The distance between molecules in an ideal gas $V/N^{1/3}$ is found from

$$PV = nRT = nRT \cdot N_A/N_A = NkT \rightarrow V/N^{1/3} = kT/P^{1/3}$$

and equating this to λ above, $kT/P^{1/3} = \frac{h}{3mkT^{1/2}}$

$$\frac{kT}{P} = \frac{h^3}{3mkT^{3/2}} \text{ and solving for } T, \text{ yields: } T^{5/2} = \frac{P}{k} \frac{h^3}{3mk^{3/2}}$$

$$T = \left[\frac{Ph^3}{k \cdot 3mk^{3/2}} \right]^{2/5} = \left[\frac{101kPa \cdot 6.63 \times 10^{-34} J \cdot s^3}{3 \cdot 2 \times 1.67 \times 10^{-27} kg \cdot 1.38 \times 10^{-23} J/K^{5/2}} \right]^{2/5} = 4.4K$$

$$8-26. \quad T_C = \frac{h^2}{2mk} \left[\frac{N}{2\pi \cdot 2.315 V} \right]^{2/3} \quad (\text{Equation 8-48})$$

The density of liquid Ne is 1.207 g/cm^3 , so

$$\frac{N}{V} = \frac{1.207 \text{ g/cm}^3 \cdot 6.022 \times 10^{23} \text{ molecules/mol} \cdot 10^6 \text{ cm}^3/\text{m}^3}{20.18 \text{ g/mol}} = 3.601 \times 10^{28} / \text{m}^3$$

$$T = \frac{6.626 \times 10^{-34} J \cdot s^2}{2 \cdot 20u \times 1.66 \times 10^{-27} kg/u \cdot .381 \times 10^{-23} J/K} \left[\frac{3.601 \times 10^{28} \text{ m}^3}{2\pi \cdot 2.315} \right]^{2/3} = 0.895K$$

Thus, T_C at which ^{20}Ne would become a superfluid is much lower than its freezing temperature of $24.5K$.

8-28. $\langle E \rangle = \frac{hf}{e^{hf/kT} - 1}$ (Equation 8-60)

(a) For $T = 10hf/k$, $hf = kT/10 \rightarrow \langle E \rangle = \frac{hf}{e^{1/10} - 1} = \frac{kT/10}{0.1051} = 0.951kT$

(b) For $T = hf/k$, $hf = kT \rightarrow \langle E \rangle = \frac{hf}{e^1 - 1} = \frac{kT}{1.718} = 0.582kT$

(c) For $T = 0.1hf/k$, $hf = 10kT \rightarrow \langle E \rangle = \frac{hf}{e^{10} - 1} = \frac{10kT}{2.20 \times 10^4} = 4.54 \times 10^{-4}kT$

According to equipartition $\langle E \rangle = kT$ in each case.

8-29. $C_V = 3N_A k \left(\frac{hf}{kT} \right)^2 \frac{e^{hf/kT}}{e^{hf/kT} - 1}^2$ As $T \rightarrow \infty$, hf/kT gets small and

$$e^{hf/kT} \approx 1 + hf/kT + \dots$$

$$C_V = 3N_A k \left(\frac{hf}{kT} \right)^2 \frac{1 + hf/kT + \dots}{hf/kT}^2 \approx 3N_A k = 3N_A R/N_A = 3R$$

The rule of Dulong and Petit.

8-31. $C_V = 3R \left(\frac{hf}{kT} \right)^2 \frac{e^{hf/kT}}{e^{hf/kT} - 1}^2$ (Equation 8-62)

At the Einstein temperature $T_E = hf/k$,

$$C_V = 3R \frac{e^1}{e^1 - 1}^2 = 3R \frac{0.9207}{0.9207} = 3 \frac{8.31 J/K \cdot mol}{0.9207}$$

$$= 22.95 K / K \cdot mol = 5.48 cal / K \cdot mol$$

8-35. Approximating the nuclear potential with an infinite square well and ignoring the Coulomb repulsion of the protons, the energy levels for both protons and neutrons are given by $E_n = n^2 h^2 / 8mL^2$ and six levels will be occupied in ^{22}Ne , five levels with 10 protons and six levels with 12 neutrons.

$$E_F \text{ protons} = \frac{5^2 \cdot 1240 \text{ MeV} \cdot \text{fm}^2}{8 \cdot 1.0078 u \times 931.5 \text{ MeV} / u \cdot 3.15 \text{ fm}^2} = 516 \text{ MeV}$$

$$E_F \text{ neutrons} = \frac{6^2 \cdot 1240 \text{ MeV} \cdot \text{fm}^2}{8 \cdot 1.0087 u \times 931.5 \text{ MeV} / u \cdot 3.15 \text{ fm}^2} = 742 \text{ MeV}$$

$$\langle E \rangle \text{ protons} = 3/5 E_F = 310 \text{ MeV}$$

$$\langle E \rangle \text{ neutrons} = 3/5 E_F = 445 \text{ MeV}$$

As we will discover in Chapter 11, these estimates are nearly an order of magnitude too large. The number of particles is not a large sample.

8-36. $E_1 = h^2 / 8mL^2$. All 10 bosons can be in this level, so $E_1 \text{ total} = 10h^2 / 8mL^2$.

8-39. For a one-dimensional well approximation, $E_n = n^2 h^2 / 8mL^2$. At the Fermi level E_F , $n=N/2$, where $N =$ number of electrons.

$$E_F = \frac{N/2^2 h^2}{8mL^2} = \frac{h^2}{32m} \left(\frac{N}{L} \right)^2 \quad \text{where } N/L = \text{number of electrons/unit length,}$$

i.e., the density of electrons. Assuming 1 free electron/Au atom,

$$\frac{N}{L} = \left[\frac{N_A \rho}{M} \right]^{1/3} = \left[\frac{6.02 \times 10^{23} \text{ electrons/mol} \cdot 19.32 \text{ g/cm}^3 \cdot 10^2 \text{ cm/m}^3}{197 \text{ g/mol}} \right]^{1/3} = 3.81 \times 10^9 \text{ m}^{-1}$$

$$E_F = \frac{6.63 \times 10^{-34} \text{ J}\cdot\text{s}^2 \cdot 3.81 \times 10^9 \text{ m}^{-1}^2}{32 \cdot 9.11 \times 10^{-31} \text{ kg} \cdot 1.602 \times 10^{-19} \text{ J/eV}} = 1.37 \text{ eV}$$

This is the energy of an electron in the Fermi level above the bottom of the well. Adding the work function to such an electron just removes it from the metal, so the well is $1.37 \text{ eV} + 4.8 \text{ eV} = 6.2 \text{ eV}$ deep.