

Problem 1

$$E_n = E_1 n^2, \text{ with } E_1 = \frac{\hbar^2 \pi^2}{2m_e L^2}$$

$$hf = E_n - E_{n-1} = E_1 (n^2 - (n-1)^2) = E_1 (2n-1) = \frac{\hbar^2 \pi^2}{2m_e L^2} (2n-1)$$

$$\Rightarrow f = \frac{\hbar^2 \pi^2}{2m_e L^2} (2n-1) = \frac{k \cdot \hbar \cdot \pi^2}{2\pi \cdot 2m_e L^2} (2n-1) \Rightarrow \boxed{f = \frac{\pi \hbar}{4m_e L^2} (2n-1)} \quad (a)$$

(b) $E = \frac{p^2}{2m_e} \Rightarrow p^2 = 2m_e E \Rightarrow \text{in } n\text{-th state, } p^2 = 2m_e E_n \Rightarrow$

$$p^2 = 2m_e E_1 n^2 = \frac{\hbar^2 \pi^2}{L^2} n^2 \Rightarrow p = \frac{\pi \hbar}{L} n$$

Now $p = m_e v \Rightarrow v = \frac{\pi \hbar}{m_e L} n$; electron is moving back and forth with speed $v \Rightarrow$ travels distance $2L$ in time $t = \frac{2L}{v} = \frac{2L^2 m_e}{\pi \hbar n}$

\Rightarrow the classical frequency of oscillation is

$$\boxed{f = \frac{1}{t} = \frac{\pi \hbar}{2m_e L^2} n}$$

(c) For large n , $2n-1 \approx 2n$, so the quantum frequency is

$$f = \frac{\pi \hbar}{4m_e L^2} (2n-1) \approx \frac{\pi \hbar}{2m_e L^2} n$$

\Rightarrow agrees with classical frequency for large n , as expected from the correspondence principle.

Problem 2

$$E_n = \hbar\omega(n + \frac{1}{2}) ; \text{ first excited state} \Rightarrow n = 1 \Rightarrow n + \frac{1}{2} = \frac{3}{2} \Rightarrow$$

$$\Rightarrow E_1 = \frac{3}{2}\hbar\omega = \frac{1}{2}m_e\omega^2 A_1^2, \text{ where } A_1 \text{ is the classical amplitude}$$

$$\Rightarrow \hbar^2 E_1 = \frac{1}{2}m_e(\hbar\omega)^2 A_1^2 = \frac{1}{2}m_e\left(\frac{2}{3}\right)^2 \cdot \left(\frac{3}{2}\hbar\omega\right)^2 \cdot A_1^2 \Rightarrow$$

$$\Rightarrow \hbar^2 E_1 = \frac{2}{9}m_e E_1 / A_1^2 \Rightarrow E_1 = \frac{9\hbar^2}{2m_e A_1^2} \Rightarrow$$

$$\Rightarrow E_1 = 7.62 \text{ eV} \text{ Å}^2 \cdot \frac{9}{2} \cdot \frac{1}{25 \text{ Å}^2} \Rightarrow \boxed{E_1 = 1.37 \text{ eV}}$$

$$(b) \text{ Probability} \Rightarrow P(x) = |\psi(x)|^2 = C^2 x^2 e^{-\frac{m_e\omega}{\hbar^2} x^2}$$

$$\text{Find max: } P'(x_m) = 0 = 2x_m - x_m^2 \cdot \frac{m_e\omega}{\hbar^2} \cdot 2x_m \Rightarrow \boxed{x_m = \left(\frac{\hbar}{m_e\omega}\right)^{1/2}}$$

(c) The classical amplitude is:

$$\frac{3}{2}\hbar\omega = \frac{1}{2}m_e\omega^2 A_1^2 \Rightarrow A_1^2 = \frac{x}{m_e\omega} \times \frac{3}{2}\hbar\omega = \frac{3\hbar}{m_e\omega}$$

$$\Rightarrow A = \sqrt{3} \left(\frac{\hbar}{m_e\omega}\right)^{1/2} \Rightarrow x_m = \frac{A}{\sqrt{3}}, A = \sqrt{3}x_m$$

Ratio of probabilities: since $P(x) \propto x^2 e^{-x^2/x_m^2} \Rightarrow$

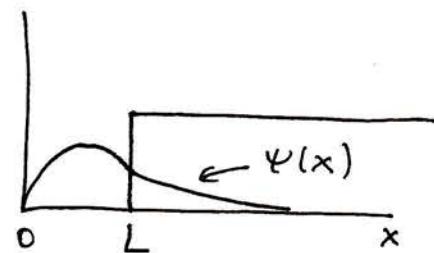
$$\frac{P(x_m)}{P(A_1)} = \frac{e^{-x_m^2/x_m^2}}{e^{-3x_m^2/x_m^2}} \cdot \frac{x_m^2}{3x_m^2} = \frac{e^3}{3} = 6.7$$

\Rightarrow electron is 6.7 times more likely to be at x_m than at the classical amplitude

Problem 3

(i) $\Psi(x) = A \sin kx + B \cos kx$

(ii) $\Psi(x) = C e^{-\alpha x} + D e^{\alpha x}$



(a) Since $V(x) = \infty$ for $x < 0 \Rightarrow \Psi(x) = 0$ for $x < 0 \Rightarrow$ by continuity, $\Psi(x=0) = 0 \Rightarrow [B=0]$

Since wavefunction has to be normalizable and $e^{\alpha x} \xrightarrow{x \rightarrow \infty} \infty \Rightarrow [D=0]$

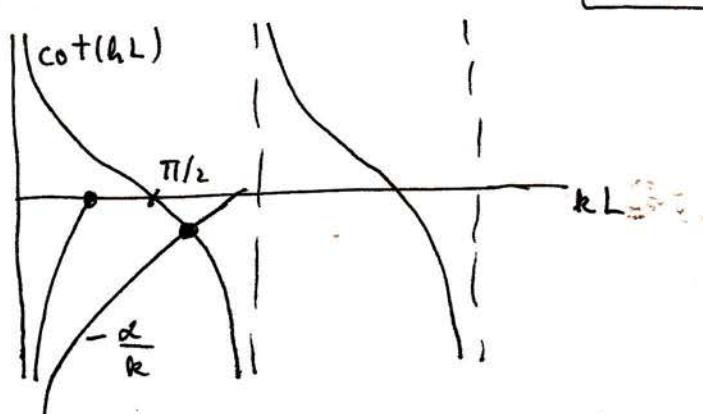
(b) Since wavefunction satisfies Schrödinger eq. in all regions

$$k = \left(\frac{2me}{\hbar^2} E \right)^{1/2} = \left(\frac{2 \cdot 3}{7.62} \right)^{1/2} \text{Å}^{-1} = [0.887 \text{Å}^{-1}]; \alpha = \sqrt{\frac{2m}{\hbar^2} (V_0 - E)} = 1.02 \text{Å}^{-1}$$

(c) See graph of $\Psi(x)$ in figure above.

(d) $\Psi(x) = A \sin kx$ for $0 < x < L$, $\Psi(x) = C e^{-\alpha x}$ for $x > L \Rightarrow$

Continuity $\Rightarrow A \sin kL = C e^{-\alpha L} \Rightarrow \cot(kL) = -\frac{\alpha}{k}$
 cont. of Ψ' $\Rightarrow kA \cos kL = -\alpha C e^{-\alpha L} \Rightarrow \cot(kL) = -\frac{\alpha}{k}$



The figure shows two examples of possible functions $-\alpha/k$.
 One of them intersects $\cot(kL) \Rightarrow$ there is a solution
 the other one does not \Rightarrow no solution.
 \therefore there is no solution if V_0 is very small