

Physics 140B, Winter 2010
Midterm exam
Feb 16, 11:00 a.m. – 12:20 p.m.

Notes: (i) Attempt as many questions as you can.
(ii) All questions may not be of equal value!

1. (a) Show that the Debye temperature of a solid is given by the formula

$$\Theta_D = (hc/k) (3N / 4\pi V)^{1/3},$$

where the various symbols have their usual meanings; note that the symbol c here stands for the speed of sound waves --- both longitudinal and transverse.

(b) For copper (atomic weight 63.6), the mass density ρ is $8.95 \times 10^3 \text{ kg.m}^{-3}$ while the speed of sound c is $2.24 \times 10^3 \text{ m.s}^{-1}$. Using these data, determine Θ_D for copper.

2. Consider a gas of N identical atoms trapped in a magnetic field such that they behave like indistinguishable simple harmonic oscillators obeying Bose-Einstein statistics. The energy of any one of those atoms is given by the expression

$$\varepsilon = p^2 / 2m + (1/2)Kr^2,$$

so that they all vibrate with a common frequency

$$\nu = (1/2\pi) (K/m)^{1/2}.$$

(a) Show that, at low temperatures, this gas undergoes the phenomenon of Bose-Einstein condensation, with

$$T_B \sim (h\nu/k) N^{1/3}.$$

(b) Next, determine the manner in which the condensate fraction N_0/N and the specific heat C_v of the gas vary with T --- when $T < T_B$.

3. The Helmholtz free energy of a non-relativistic Fermi gas at low temperatures is given by the expression

$$F = (3/5)N \varepsilon_F \left[1 - (5\pi^2/12) (kT / \varepsilon_F)^2 + \dots \right],$$

where ε_F is the Fermi energy of the gas; note that $\varepsilon_F \propto (N/V)^{2/3}$.

(a) Using this expression for F , derive the corresponding expression for the chemical potential μ of the gas.

4. Consider an electron gas with particle density n . Determine the *numerical* value of n for which the Fermi momentum p_F of the gas is equal to mc , where m is the rest-mass of the electron and c the speed of light.

What is the corresponding value of the Fermi energy ε_F (in units of mc^2) and the Fermi velocity v_F (in units of c)?

5. Using Maxwell's relations, show that for a magnetic system

$$\left(\frac{\partial C_H}{\partial H} \right)_T = T \left(\frac{\partial^2 M}{\partial T^2} \right)_H.$$

Apply this result to a paramagnetic material *in the Curie regime* and show that, in this case,

$$C_H = C H^2 / T^2 + f(T),$$

where C is the Curie constant and $f(T)$ an unknown function of T .

1. (a) $\Theta_D = h\nu_{\max}/k$, where ν_{\max} is determined by the condition

$$\int_0^{\nu_{\max}} 4\pi V \left(\frac{1}{c_l^3} + \frac{2}{c_t^3} \right) \nu^2 d\nu = 3N.$$

Dropping the distinction between c_l & c_t , we get

$$4\pi V \frac{3}{c^3} \frac{\nu_{\max}^3}{3} = 3N. \quad \therefore \nu_{\max} = c \left(\frac{3N}{4\pi V} \right)^{1/3}.$$

It follows that

$$\textcircled{4} \quad \Theta_D = \frac{hc}{k} \left(\frac{3N}{4\pi V} \right)^{1/3} \checkmark$$

(b) Since $\rho = Nm_{cu}/V$, we get

$$\begin{aligned} \Theta_D &= \frac{hc}{k} \left(\frac{3\rho}{4\pi m_{cu}} \right)^{1/3} \\ &= \frac{6.626 \times 10^{-34} \text{ Js} \cdot 2.24 \times 10^3 \text{ m s}^{-1}}{1.381 \times 10^{-23} \text{ JK}^{-1}} \left(\frac{3 \times 8.95 \times 10^3 \text{ kg m}^{-3}}{4\pi \times 63.6 \times 1.661 \times 10^{-27} \text{ kg}} \right)^{1/3} \end{aligned}$$

$$\textcircled{4} \quad = 293 \text{ K.} \checkmark$$

2. (a) In the region of BEC, $\mu = \epsilon_{\min} \approx 0$. So, the expectation value of the total # of atoms in all the excited states is

$$N_{\text{exc}} = \sum_{\epsilon > \epsilon_{\min}} \overline{N}_{\epsilon} \approx \int_0^{\infty} \frac{g(\epsilon) d\epsilon}{e^{\epsilon/kT} - 1}$$

$$\approx \int_0^{\infty} \int_0^{\infty} \frac{4\pi p^2 dp \cdot 4\pi r^2 dr / h^3}{e^{p^2/2mkT + Kr^2/2kT} - 1}$$

Set $p = \sqrt{2mkT} u$ & $r = \sqrt{2kT/K} v$, to get

$$N_{\text{exc}} = 128 \pi^2 (m/K)^{3/2} (kT/h)^3 \cdot I,$$

where $I = \int_0^{\infty} \int_0^{\infty} \frac{u^2 v^2 du dv}{e^{u^2+v^2} - 1}$, a numerical constant.

It follows that

$$N_{\text{exc}} = \text{const.} (kT/hv)^3$$

Now, T_B is determined by the condition: $N_{\text{exc}} = N$, with the result that

(6) $T_B = \text{const.} (hv/k) \cdot N^{1/3}$ ✓

(b) For $T < T_B$, $N_{\text{exc}} \propto T^3$ and, at $T = T_B$, $N_{\text{exc}} = N$. It follows that, for $T < T_B$,

$$(3) \quad N_{exc}/N = (T/T_B)^3$$

and, hence,

$$N_0/N = 1 - (T/T_B)^3. \quad \checkmark$$

As for C_V , we note ^{that} the mean thermal energy carried by an atom in an excited state is $O(kT)$. The total thermal energy of the gas is, therefore, $\propto T^4$ and, hence, its $C_V \propto T^3$. \checkmark

(3)

3. Since $dF = -SdT - PdV + \mu dN$, therefore $\mu = \left(\frac{\partial F}{\partial N} \right)_{T,V}$.

The given expression for F is

$$F = \frac{3}{5} N \varepsilon_F - \frac{\pi^2}{4} N \frac{k^2 T^2}{\varepsilon_F} + \dots$$

Since $\varepsilon_F \propto (N/V)^{2/3}$, the first term of F is $\propto N^{5/3}$ and the second term is $\propto N^{1/3}$; other factors are functions of T & V only, so they remain unaffected by differentiation w.r.t. N . We thus get

$$\mu = \frac{5}{3N} \cdot \frac{3}{5} N \varepsilon_F - \frac{1}{3N} \cdot \frac{\pi^2}{4} N \frac{k^2 T^2}{\varepsilon_F} + \dots$$

(8)

$$= \varepsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{kT}{\varepsilon_F} \right)^2 + \dots \right] \checkmark$$

4. The Fermi momentum p_F is given by the condition

$$\int_0^{p_F} 2 \cdot \frac{V 4\pi p^2 dp}{h^3} = N, \text{ i.e. } p_F = h \left(\frac{3N}{8\pi V} \right)^{1/3}.$$

Equating p_F with mc , we get

$$(4) \quad n = \frac{N}{V} = \frac{8\pi}{3} \left(\frac{mc}{h} \right)^3 = 5.84 \times 10^{35} \text{ m}^{-3} \checkmark$$

The corresponding values of ϵ_F/mc^2 and v_F/c are given by

$$(2) \quad \epsilon_F = c \sqrt{p_F^2 + m^2 c^2} - mc^2 = (\sqrt{2} - 1) mc^2, \therefore \frac{\epsilon_F}{mc^2} = \sqrt{2} - 1, \text{ and } \checkmark$$

$$(2) \quad v_F = \frac{p_F}{m_{rel}} = \frac{p_F c^2}{\epsilon_{total}} = \frac{mc \cdot c^2}{\sqrt{2} mc^2} = \frac{c}{\sqrt{2}}, \therefore \frac{v_F}{c} = \frac{1}{\sqrt{2}} \checkmark$$

$$5. \left(\frac{\partial C_H}{\partial H} \right)_T = \left[\frac{\partial}{\partial H} \left\{ T \left(\frac{\partial S}{\partial T} \right)_H \right\} \right]_T$$

$$= T \frac{\partial^2 S}{\partial H \partial T} = T \frac{\partial^2 S}{\partial T \partial H}$$

Using one of Maxwell's relations, we get

$$(4) \quad T \left[\frac{\partial}{\partial T} \left(\frac{\partial M}{\partial T} \right)_H \right]_H = T \left(\frac{\partial^2 M}{\partial T^2} \right)_H \quad \checkmark$$

For a paramagnetic material in the Curie regime,

$$M = \frac{CH}{T}. \quad \text{Therefore, } \left(\frac{\partial^2 M}{\partial T^2} \right)_H = \frac{2CH}{T^3}.$$

It follows that

$$\left(\frac{\partial C_H}{\partial H} \right)_T = \frac{2CH}{T^2}.$$

Integrating over H , keeping T const., we get

$$(4) \quad C_H = \frac{CH^2}{T^2} + \text{a const. independent of } H, \quad \checkmark$$

but possibly a function of T .

Hence the result.