

Phy 212 C. HW 1

1. Consider a bosonic system whose single particle ^{basis} is denoted as $\psi_\alpha, \psi_\beta, \dots$. The particle number in each state is denoted n_α, n_β, \dots . The many-body eigenstates in the particle-number occupation representation is denoted $|n_\alpha n_\beta \dots\rangle$. We use a_α^\dagger and a_α as the creation and annihilation operators for the state α , which satisfies $[a_\alpha, a_\beta^\dagger] = i\delta_{\alpha\beta}$.

a) Prove $a_\alpha^\dagger |n_\alpha n_\beta \dots\rangle = \sqrt{n_\alpha + 1} |n_\alpha + 1, n_\beta, \dots\rangle$

$a_\alpha |n_\alpha n_\beta \dots\rangle = \sqrt{n_\alpha} |n_\alpha - 1, n_\beta, \dots\rangle$

b) let $\hat{n}_\alpha = a_\alpha^\dagger a_\alpha$, prove $\hat{n}_\alpha |n_\alpha n_\beta \dots\rangle = n_\alpha |n_\alpha n_\beta \dots\rangle$

c) Prove $|n_\alpha n_\beta \dots\rangle = \frac{1}{\sqrt{n_\alpha! n_\beta! \dots}} (a_\alpha^\dagger)^{n_\alpha} (a_\beta^\dagger)^{n_\beta} \dots |0_\alpha 0_\beta \dots\rangle$

Schwinger Boson:

2.

② Let us consider two independent harmonic oscillators to form a system. Let us use n_1, n_2 to denote the number of phonons, and $a_1^\dagger, a_1, a_2^\dagger, a_2$ for the creation / annihilation operators, $n_1 = a_1^\dagger a_1, n_2 = a_2^\dagger a_2$ are number operators. In the phonon number representation, the phonon eigenstates basis are denoted as $|n_1, n_2\rangle$. We write $a = \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}$, $a^\dagger = (a_1^\dagger, a_2^\dagger)$ in the matrix form

and define $\vec{J} = \frac{1}{2} a_\alpha^\dagger \vec{\sigma}_{\alpha\beta} a_\beta$, where $\vec{\sigma}$ is the usual Pauli matrix.

1) define $J_\pm = J_x \pm iJ_y$, show $J_+ = a_1^\dagger a_2, J_- = a_2^\dagger a_1$

2) Prove $[J_z, J_\pm] = \pm J_\pm$, and $[J_+, J_-] = 2J_z$, which are the angular momentum algebra.

3) Calculate $J^2 = J_x^2 + J_y^2 + J_z^2$, and express it in terms of n_1 and n_2 .

4) Show that the state $|n_1, n_2\rangle$ is an eigenstate of J^2, J_z with

$$j = \frac{1}{2}(n_1 + n_2) \text{ and } j_z = \frac{1}{2}(n_1 - n_2),$$

5) Prove $J_\pm |j, m\rangle = \sqrt{(j \mp m)(j \pm m + 1)} |j, m \pm 1\rangle$ by using the Schwinger identity in angular momentum theory

bosonic representation defined above.

6) Define operators $K_+ = a_1^\dagger a_2^\dagger, K_- = a_1 a_2$, calculate their effects when acting on $|n_1, n_2\rangle$, and then interpret in the angular momentum representation.

3: Explicitly verify the equation of continuity

$$\frac{\partial}{\partial t} \rho(\mathbf{r}, t) + \nabla \cdot \vec{j}(\mathbf{r}, t) = 0 \text{ for a free Fermion gas.}$$

The Hamiltonian $H_0 = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} C_{\mathbf{k}\sigma}^\dagger C_{\mathbf{k}\sigma}$. The density operator is

defined as $\rho(\mathbf{r}, t) = e^{iHt} \rho(\mathbf{r}) e^{-iHt}$, where $\rho(\mathbf{r}) = \psi_\sigma^\dagger(\mathbf{r}) \psi_\sigma(\mathbf{r}) = \frac{1}{V} \sum_{\mathbf{k}\sigma} e^{-i\mathbf{q}\cdot\mathbf{r}} C_{\mathbf{k}+\mathbf{q}\sigma}^\dagger C_{\mathbf{k}\sigma}$.

The current operator $\mathbf{j}(\mathbf{r}, t) = e^{iHt} \mathbf{j}(\mathbf{r}) e^{-iHt}$ and

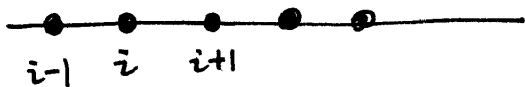
$$\mathbf{j}(\mathbf{r}, t) = \frac{-i\hbar}{m} [\psi_\sigma^\dagger(\mathbf{r}) \nabla \psi_\sigma(\mathbf{r}) - (\nabla \psi_\sigma^\dagger(\mathbf{r})) \psi_\sigma(\mathbf{r})]$$

$$= \frac{1}{V} \sum_{\mathbf{k}\sigma} e^{-i\mathbf{q}\cdot\mathbf{r}} \frac{(\vec{k} + \frac{1}{2}\vec{q})}{m} C_{\mathbf{k}+\mathbf{q}\sigma}^\dagger C_{\mathbf{k}\sigma}.$$

4: Do the Cooper problem again as in the lecture notes and find the gap value. (If you had done this problem 60 years before, you would be awarded the Nobel prize).

5: Tight binding model. Consider in solid, we denote C_i^\dagger, C_i for the creation/annihilation operators for electrons on site i .

The hopping Hamiltonian is



$$H = -t \sum_{i\sigma} C_{i\sigma}^\dagger C_{i+1\sigma}.$$

By introducing the discrete Fourier transform to diagonalize the spectra $H = \sum_{\mathbf{k}\sigma} \epsilon(\mathbf{k}) C_{\mathbf{k}\sigma}^\dagger C_{\mathbf{k}\sigma}$, where $C_{\mathbf{k}} = \frac{1}{\sqrt{N}} \sum_{i \rightarrow \text{number of sites}} e^{i\mathbf{k}\cdot\vec{r}_i} C_i$.