

$$= \int_0^u \gamma^3 m u^2 + m c^2 \sqrt{1 - \frac{u^2}{c^2}} \Big|_0^u$$

$$= \gamma^3 m u^2 + m c^2 / \gamma - m c^2$$

$$\gamma T = \gamma^3 m u^2 + m c^2 - \gamma m c^2$$

$$= \frac{m u^2}{1 - u^2/c^2} + m c^2 - \gamma m c^2$$

$$= \frac{m u^2 + m c^2 - m u^2}{1 - u^2/c^2} - \gamma m c^2$$

$$= \gamma^2 m c^2 - \gamma m c^2$$

$$\therefore T = \gamma m c^2 - m c^2$$

$$T = (\gamma - 1) m c^2$$

For small u , $(1 - \frac{u^2}{c^2})^{-1/2} \approx 1 + \frac{1}{2} \frac{u^2}{c^2}$

and $T \rightarrow \frac{1}{2} m u^2$

So, we recover Newtonian formula.

The quantity

$$E_0 \equiv m c^2$$

is the rest energy.

We then have

$$\gamma m c^2 = T + E_0$$

We identify the quantity γmc^2 with the total energy, since it is the sum of kinetic and rest mass.

$$\boxed{E \equiv \gamma mc^2 = T + E_0} \quad \star$$

Note as $v \rightarrow 0$, $\gamma \rightarrow 1$ and $E \rightarrow E_0$.
The equation $E_0 = mc^2$ says that mass has an energy equivalent.

Thus, rather than thinking of separate conservation laws for mass and energy in relativity we have conservation of mass-energy.

Example 14.8 Binding energy of D.
(see TM)

We can further manipulate \star .
we have

$$\begin{aligned} p &= \gamma m u \\ p^2 c^2 &= \gamma^2 m^2 u^2 c^2 \\ &= \gamma^2 m^2 c^4 \left(\frac{u^2}{c^2} \right) \end{aligned}$$

but $\frac{u^2}{c^2} = 1 - \frac{1}{\gamma^2}$, so

$$\begin{aligned} p^2 c^2 &= \gamma^2 m^2 c^4 \left(1 - \frac{1}{\gamma^2} \right) \\ &= \gamma^2 m^2 c^4 - m^2 c^4 \end{aligned}$$

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Example : Solar fusion

The Sun shines by virtue of energy released when 4 H nuclei (p) are converted to 1 He nucleus (α) by nuclear fusion. The energy released per reaction is simply calculated if we know the masses:

$$m_p = 1.0072766 \text{ u}$$

$$m_\alpha = 4.000408 \text{ u}$$

$$\begin{aligned} \Delta E &= 4m_p c^2 - m_\alpha c^2 \\ &= (0.028698 \text{ u}) c^2 \end{aligned}$$

$$(1 \text{ u}) c^2 = 931.48 \text{ MeV}$$

$$\therefore \Delta E = 26.73 \text{ MeV}$$

This energy comes out in the form of 2 γ 's and 2 H_e^+ per reaction. Only the γ 's effectively heat the Sun.

but $E_0^2 = m^2 c^4$, $E^2 = \gamma^2 m^2 c^4$

so

$$p^2 c^2 = E^2 - E_0^2$$

or

$$\boxed{E^2 = p^2 c^2 + E_0^2}$$

This relates a particle's (body's) momentum, rest energy, and total energy. Since a photon has no mass, $E_0 = 0$, hence

$$E = pc \quad (\text{photon})$$

Since

$$E = h\nu$$

we have

$$\boxed{p = \frac{h\nu}{c}}$$

12. Spacetime and Four-Vectors

Because space and time become mixed in a LT, it is convenient to think of a 4-dimensional continuum called spacetime, first introduced by Minkowski.

An event in ST has coordinates

$$\text{Event: } (x_1, x_2, x_3, ct)$$

The ST interval between two events is invariant under LT

$$\Delta S_{12} = (\Delta x_1^2 + \Delta x_2^2 + \Delta x_3^2 - c^2 \Delta t^2)^{1/2}$$

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$$\Delta S'_{12} = (\Delta X_1'^2 + \Delta X_2'^2 + \Delta X_3'^2 - c^2 \Delta t'^2)^{1/2}$$

$$\Delta S'_{12} = \Delta S$$

If we introduce the variable $X_4 = ict$ ($c = \sqrt{-1}$), then we can write the ST interval as

$$\Delta S = \left(\sum_{\mu=1}^4 \Delta X_{\mu}^2 \right)^{1/2}$$

$$\Delta S' = \left(\sum_{\mu=1}^4 \Delta X_{\mu}'^2 \right)^{1/2} = \Delta S$$

In the 4-dimensional space (x_1, x_2, x_3, x_4) , ΔS is just the distance between 2 points. Since $\Delta S = \Delta S'$, the LT is distance preserving, it corresponds to a rotation in Minkowski space.

⇒ The LT is an orthogonal transformation in Minkowski space. In analogy with 3D rotations, we can write

$$X_{\mu}' = \sum_{\nu} \lambda_{\mu\nu} X_{\nu}$$

rotation matrix
in 4-D Minkowski
space

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In the case of, K' moving along X_1 , with relative velocity $v = \beta c$, we have from inspection

$$[\Lambda] = \begin{pmatrix} \gamma & 0 & 0 & i\beta\gamma \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\beta\gamma & 0 & 0 & \gamma \end{pmatrix}$$

Check

$$\begin{pmatrix} x_1' \\ x_2' \\ x_3' \\ ict' \end{pmatrix} = \begin{pmatrix} \gamma & 0 & 0 & i\beta\gamma \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -i\beta\gamma & 0 & 0 & \gamma \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ ict \end{pmatrix}$$

$$\begin{aligned} x_1' &= \gamma x_1 + (i\beta\gamma)(ict) \\ &= \gamma(x_1 - vt) \end{aligned}$$

$$x_2' = x_2$$

$$x_3' = x_3$$

$$ict' = -i\beta\gamma x_1 + \gamma(ict)$$

$$\begin{aligned} t' &= \gamma t - \frac{\beta\gamma x_1}{c} \\ &= \gamma \left(t - \frac{x_1 v}{c^2} \right) \end{aligned}$$

Any quantity that transforms according to

$$A'_\mu = \sum_\nu \Lambda_{\mu\nu} A_\nu$$

is called a 4-vector. Obviously, the coordinates of an event in Minkowski space (MS) is a 4-vector

$$\begin{aligned} X &= (x_1, x_2, x_3, ict) \\ &= (\vec{x}, ict) \end{aligned}$$

Relativistic momentum and energy

Armed with the above, we now show that \vec{p}_{rel}, E_{rel} are elements of a 4-vector. First, a little jargon.

Any scalar quantity that is invariant under a LT is called a 4-scalar. Clearly, the ST interval is a 4-scalar.

$$ds = \sqrt{\sum_\mu dx_\mu^2}$$

Now, in a frame comoving with the clock ($dx_1 = dx_2 = dx_3 = 0; dt = d\tau$)

$$ds = \sqrt{0+0+0-c^2d\tau^2}$$

$$d\tau = \frac{i}{c} ds$$

$\therefore d\tau$ is also a 4-scalar.

If X is a 4-vector, dX is a 4-vector and $\frac{dX}{dt} = V$ is a 4-vector; the 4-velocity

$$\begin{aligned} V &= \frac{dX}{dt} = \left(\frac{d\vec{x}}{dt}, ic \frac{dt}{dt} \right) \\ &= \left(\frac{d\vec{x}}{dt}, ic \right) \frac{dt}{d\tau} \\ &= \gamma (\vec{u}, ic) \end{aligned}$$

where $\gamma = \frac{1}{\sqrt{1-\beta^2}}$, $\beta = |\vec{u}|/c$

Since rest mass is a 4-scalar, the 4-momentum TP is also a 4-vector.

$$TP \equiv mV$$

$$TP = \left(\frac{m\vec{u}}{\sqrt{1-\beta^2}}, iP_4 \right)$$

where $P_4 = \frac{mc}{\sqrt{1-\beta^2}}$

The first 3 components are our relativistic momentum \vec{p}_{rel} , while the 4th component is related to E via

$$p_4 = \partial \text{mc} = \frac{E}{c}$$

Therefore, the 4-vector momentum can be written

$$\mathbb{P} = (\vec{p}_{\text{rel}}, i \frac{E}{c})$$

Since \mathbb{P} is a 4-vector, it transforms according to

$$P'_\mu = \sum_\nu \lambda_{\mu\nu} P_\nu$$

Explicitly, apply matrix for $K \rightarrow K'$

$$\begin{aligned} p'_1 &= \frac{p_1 - (v/c^2)E}{\sqrt{1-\beta^2}} \\ p'_2 &= p_2 \\ p'_3 &= p_3 \\ E' &= \frac{E - vp}{\sqrt{1-\beta^2}} \end{aligned}$$

FIN